

HW 7

① Beats at 3 Hz (3 times per second).

② (a) $E = E_L + E_R = E_0 \left(\sin(kx + \omega t) + p \sin(kx - \omega t) \right)$

→ Choose x -axis going to right. Thus E_L moves in $-x$ direction & E_R in $+x$:

$$E = E_0 \left[\sin(kx + \omega t) + p \sin(kx - \omega t) \right]$$

$$= E_0 \left[\sin(kx) \cos(\omega t) + \sin(\omega t) \cos(kx) + p \sin(kx) \cos(\omega t) + p \sin(-\omega t) \cos(kx) \right]$$

$$= \sin(kx) \cos(\omega t) \cdot E_0 p + \underline{E_0 p \sin(kx) \cos(\omega t)} - \underline{E_0 p \sin(kx) \cos(\omega t)} + E_0 \sin(kx) \cos(\omega t) + E_0 \sin(\omega t) \cos(kx) - E_0 p \sin(\omega t) \cos(kx)$$

$$= 2E_0 p \sin(kx) \cos(\omega t) + E_0 \sin(kx + \omega t) - E_0 p \sin(kx + \omega t)$$

$$E = 2E_0 p \sin(kx) \cos(\omega t) + (1-p) E \sin(kx + \omega t)$$

(b) The first term $[2E_0 \rho \sin(kx) \cos(\omega t)]$ is a standing wave, while the second is a traveling wave: $[(1-\rho)E_0 \sin(kx + \omega t)]$

When $\rho = 1$, you get 100% reflection, and a ~~static~~ standing wave is formed.

When $\rho < 1$, the reflected light is less strong than the incoming light. You will see a smaller amplitude standing wave, but there will also be an "offset" from the incoming wave.

$$(3) (a) E_{\text{arm}_1} + E_{\text{arm}_2} = E_{o1} e^{i(k\Delta x_1 + \omega t)} + E_{o2} e^{i(k\Delta x_2 + \omega t)}$$

$$\rightarrow \text{Assume } E_{o1} \approx E_{o2} \equiv E_o$$

$$\rightarrow \Delta x_1 = 2L_1 \quad \& \quad \Delta x_2 = 2L_2$$

$$E_{\text{total}} = E_o \left[e^{i(2L_1 k + \omega t)} + e^{i(2L_2 k + \omega t)} \right]$$

$$\rightarrow k = \frac{2\pi}{\lambda}$$

$$E_{\text{total}} = E_o e^{i\omega t} \left[e^{i\left[\frac{4\pi L_1}{\lambda}\right]} + e^{i\left[\frac{4\pi L_2}{\lambda}\right]} \right]$$

$$(b) I \propto E^* E = E_o^2 e^{i\omega t} e^{-i\omega t} \left[e^{i\left[\frac{4\pi L_1}{\lambda}\right]} + e^{i\left[\frac{4\pi L_2}{\lambda}\right]} \right] \\ \times \left[e^{-i\left[\frac{4\pi L_1}{\lambda}\right]} + e^{-i\left[\frac{4\pi L_2}{\lambda}\right]} \right]$$

$$I \propto E_o^2 (1) \left[e^0 + e^{-i\frac{4\pi}{\lambda}[L_2+L_1]} + e^{i\frac{4\pi}{\lambda}[L_1-L_2]} + e^0 \right]$$

$$I \propto E_o^2 \left[2 + 2 \cos\left(\frac{4\pi}{\lambda}(L_1 - L_2)\right) \right]$$

$$\rightarrow \Delta L = L_1 - L_2$$

\rightarrow a sign allows us to divide by 2 (proportionality const.)

$$I \propto E_o^2 \left(1 + \cos\left(\frac{4\pi \Delta L}{\lambda}\right) \right)$$

(3 continued)

(c) $1 + \cos(x)$ has maxima or minima every $2\pi = \Delta x$:

$$\frac{4\pi \Delta L}{\lambda_1} - \frac{4\pi \Delta L}{\lambda_2} = 2\pi$$

$$\frac{1}{\lambda_1} - \frac{1}{\lambda_2} = \frac{2\pi}{4\pi \Delta L} = \frac{1}{2\Delta L}$$

$$\frac{\lambda_2 - \lambda_1}{\lambda_1 \lambda_2} = \frac{1}{2\Delta L}$$

→ Let $\Delta\lambda = \lambda_2 - \lambda_1$ and $\lambda_1 \approx \lambda_2 \equiv \lambda$:

$$\Delta\lambda = \frac{\lambda^2}{2\Delta L}$$

(d) Resolution is ~~the~~ $\Delta\lambda$ (part c)

$$\Delta\lambda = \frac{(780.00 \text{ nm})^2}{2(1.00 - 0.10) \text{ m}} = 3.38 \times 10^{-13} \text{ m} = 0.338 \text{ pm}$$

$$\Delta\nu = \nu_1 - \nu_2 = \frac{c}{\lambda_1} - \frac{c}{\lambda_2} = \frac{c}{2\Delta L}$$

$$\Delta\nu = \text{167 MHz}$$

(3 continued)

(e) Resolution $\Delta \nu = \frac{c}{2\Delta L}$

- If you want better (i.e. smaller) resolution, ΔL should be made bigger (i.e. the interferometer arms are more unequal in length).
- ~~On~~ On the other hand, you can have the arms so unequal that the light is no longer able to interfere with itself. This is given by the coherence time & length:

$$\Delta L_c = c\Delta T_c \quad (\text{coherence time \& length})$$

Total difference in OPLs $< \Delta L_c$

$$\boxed{2(L_1 - L_2) < \Delta L_c}$$

This will limit the possible resolution you can achieve.

(4) (a) The P.D. will measure the beat between

f_{shifted} and f_0 :

$$f_{\text{shifted}} - f_0 = f_0 \left[\frac{v + v_m}{v - v_m} - 1 \right] \quad \leftarrow \text{use } v = c \text{ (vacuum)}$$

$$= f_0 \left[\frac{c + v_m - c + v_m}{c - v_m} \right]$$

$$\text{Beat} = \Delta f = f_0 \left(\frac{2v_m}{c - v_m} \right)$$

(b) $\lambda = 550 \text{ nm}$

$$f_0 = \frac{c}{\lambda} = 545 \text{ THz}$$

$$v_m = 10 \text{ cm/s}$$

$$\text{Beat} = (545 \times 10^{12} \text{ s}^{-1}) \frac{2 \cdot 10 \text{ cm/s}}{3 \times 10^8 \text{ m/s} - 0.10 \text{ m/s}}$$

$$= \boxed{364 \text{ kHz}}$$

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$$E = E_0 (1 + a \cos \omega_m t) \cos \omega_c t$$

$$= E_0 \cos \omega_c t + E_0 a \cos \omega_m t \cos \omega_c t$$

$$= E_0 \cos \omega_c t + E_0 a \left[\frac{\cos[(\omega_m + \omega_c)t] + \cos[(\omega_m - \omega_c)t]}{2} \right]$$

$$E = E_0 \cos(\omega_c t) + \frac{E_0 a}{2} (\cos[(\omega_c + \omega_m)t]) + \frac{E_0 a}{2} \cos[(\omega_c - \omega_m)t]$$

You want side bands from 20 Hz to 20,000 Hz (20 kHz),
you need this range on both sides of carrier.

$$\Delta \omega \geq 40 \text{ kHz}$$

E.C.

$$(a) \quad v_g = \frac{d\omega}{dk} = \frac{d}{d(2\pi/\lambda)} (2\pi\nu) = \frac{d\nu}{d(1/\lambda)}$$

$$\rightarrow \nu = \frac{c/n}{\lambda}$$

$$\frac{d\nu}{d(1/\lambda)} = \frac{c}{n} + \frac{c}{\lambda} \cdot \frac{d}{d(1/\lambda)} \left[\frac{1}{n} \right] \quad \leftarrow \text{Chain Rule}$$

$$(b) \text{ phase velocity} = v_{ph} = \frac{c}{n}$$

$$v_g = v_{ph} + \frac{c}{\lambda} \frac{d[1/n]}{d[1/\lambda]}$$

If $v_g < v_{ph}$, then $\frac{d[1/n]}{d[1/\lambda]}$ must be negative.

• Let $y = \frac{1}{\lambda}$. From plot, we see $n(\lambda) \approx \frac{|A|}{\lambda}$
(a good guess at the functional form). Thus, $\frac{1}{n} \approx \frac{\lambda}{|A|}$.

$$\frac{d}{d[1/\lambda]} \left[\frac{1}{n(\lambda)} \right] \approx \frac{d}{dy} \left(\frac{1}{A} \cdot \frac{1}{y} \right) = \frac{-1}{|A|y^2} = \frac{-1}{|A|} \cdot \lambda^2$$

~~Therefore~~ This term is negative! $\hookrightarrow v_g < v_{ph}$